

# 用非线性 M-C 破坏准则计算深埋洞室 围岩压力上限解

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**摘要:**在构建深埋洞室破坏机制的基础上,采用切线法将非线性 Mohr-Coulomb 破坏准则运用到极限分析上限法中,推导非线性 Mohr-Coulomb 破坏准则下深埋洞室围岩压力的上限解析解.结果表明:非线性系数、初始粘聚力是影响围岩压力的主要因素,效果非常显著,而侧压力比例系数、洞室断面尺寸以及轴向拉应力影响效果相对较小.针对深埋洞室的支护设计,采用非线性 Mohr-Coulomb 破坏准则进行计算更加科学.

**关键词:**深埋洞室;非线性 Mohr-Coulomb 破坏准则;极限分析;围岩压力;上限解

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## Upper Bound Solution to Surrounding Rock Pressure for Deep Cavity by Using Nonlinear M-C Failure Criterion

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**Abstract:** Based on constructing the failure mechanism of deep cavity, the nonlinear Mohr-Coulomb failure criterion is introduced into limit analysis by virtue of the tangential technique, and the upper bound solution of the surrounding rock pressure is finally deduced. Results show that the nonlinear coefficient and initial cohesion have greater impact on the surrounding rock pressure, while the influences of proportionality coefficient of lateral pressure, size of the cross section and axial tension stress are relatively smaller. Thus, in view of the supporting design of deep cavity, a nonlinear Mohr-Coulomb failure mechanism would be more proper.

**Keywords:** deep cavity; nonlinear M-C failure criterion; limit analysis; surrounding rock pressure; upper bound solution

深埋洞室的稳定性问题一直都是地下工程界研究的热点,而当前急需解决的难题就是如何准确获取深埋洞室破坏时的围岩压力.目前,关于求解围岩压力的解析方法主要有极限平衡法和极限分析法<sup>[1-4]</sup>.极限平衡法,应用比较广泛,但是没有考虑岩体的本构关系,而极限分析法采用正交流动法则考虑了岩体的本构关系.因此,相对极限平衡法来讲,上限理论推理严谨,尤其是不需要研究岩体弹塑性变形的全过程,而是直接关注岩体的极限破坏状态,因此计算过程简单,它是解决深埋洞室围岩压力问题的有效途径之

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一<sup>[5-9]</sup>. 近些年, 国内外一些学者采用极限分析法对深埋洞室围岩压力问题进行了相关的研究. 如冯利坡等<sup>[10]</sup>针对深埋隧道开挖面的稳定性问题, 采用极限分析法求解了不同水压下的围岩压力, 并且结合工程实例验证了该方法的正确性. Yang 等<sup>[11]</sup>将极限分析理论与变分原理相结合, 推导了深埋矩形洞室拱顶破坏围岩压力的表达式, 并且给出了相应的破坏范围. Huang 等<sup>[12]</sup>考虑孔隙水的影响, 采用极限分析法计算深埋圆形洞室的围岩压力, 并且与已有研究成果进行比较, 验证了计算结果的正确性. Zhang 等<sup>[13]</sup>采用上限法得到了深埋隧道开挖面的围岩压力, 与浅埋隧道计算结果进行比较, 变化规律的一致性验证了该方法的可行性. Qin 等<sup>[14]</sup>考虑多层土体, 采用上限法研究参数的变异性对深埋矩形洞室围岩压力的影响. 上述研究针对深埋洞室的围岩压力问题基于线性 Mohr-Coulomb 破坏准则, 而实验表明, 在软弱的围岩中, 最大主应力与最小主应力的关系是非线性关系<sup>[15]</sup>. 由于深埋洞室一般具有软弱岩层的特点, 即岩体材料服从非线性 Mohr-Coulomb 破坏准则<sup>[16-18]</sup>. 因此, 为了准确求解深埋洞室破坏时的围岩压力, 为现场设计提供科学的理论依据, 本文运用非线性 Mohr-Coulomb 破坏准则下的上限法求解深埋洞室的围岩压力.

## 1 非线性 Mohr-Coulomb 破坏准则与切线法

### 1.1 非线性 Mohr-Coulomb 破坏准则

在实际的岩土工程分析中, 非线性 Mohr-Coulomb 破坏准则得到了国内外学者们的广泛应用, 该表达式为<sup>[12]</sup>

$$\sigma_1 = q_p + M_p^* \left( \frac{\sigma_3}{q_p} \right)^\alpha. \quad (1)$$

式中:  $q_p$  为三轴压缩实验测得的土体无侧限抗压强度;  $M_p^*$ ,  $\alpha$  为经过实验测得的指标;  $\sigma_1, \sigma_3$  分别为最大和最小主应力.

### 1.2 切线法

在应力空间中, 式(2)与式(1)等价<sup>[13,14]</sup>:

$$\tau = c_0 \left( 1 + \frac{\sigma_n}{\sigma_t} \right)^{\frac{1}{m}}. \quad (2)$$

式中:  $c_0$  为初始粘聚力;  $\sigma_t$  为轴向拉应力;  $\tau$  和  $\sigma_n$  为屈服面上的剪应力和正应力;  $c_0, \sigma_t$  和  $m$  为由三轴试验确定的岩土材料参数.

式(2)绘成曲线后如图1所示. 当  $m=1$  时, 式(2)变成线性 Mohr-Coulomb 破坏准则.

切线法是采用过曲线上某一点 ( $M$  点) 的切线来代替该点的实际抗剪强度<sup>[19-21]</sup>.

切线方程为

$$\tau = c_t + \tan\varphi_t \sigma_n. \quad (3)$$

式中:  $\tan\varphi_t$  和  $c_t$  分别为切线的斜率和截距; 两者表达式各为

$$\tan\varphi_t = \frac{d\tau}{d\sigma_n} = \frac{c_0}{m\sigma_t} \left( 1 + \frac{\sigma_n}{\sigma_t} \right)^{\frac{1-m}{m}} = \frac{\tan\varphi_0}{m} \left( 1 + \frac{\sigma_n}{\sigma_t} \right)^{\frac{1-m}{m}}; \quad (4)$$

$$\sigma_n = \sigma_t \left( \frac{m\sigma_t \tan\varphi_t}{c_0} \right)^{\frac{m}{1-m}} - \sigma_t; \quad (5)$$

$$\tau = \sigma_n \tan\varphi_t + c_t = c_0 \left( \frac{m\sigma_t \tan\varphi_t}{c_0} \right)^{\frac{1}{1-m}}. \quad (6)$$

联立式(2)~式(6)可得:

$$c_t = \tau - \sigma_n \tan\varphi_t = \frac{m-1}{m} c_0 \left( \frac{m\sigma_t \tan\varphi_t}{c_0} \right)^{\frac{1}{1-m}} + \sigma_t \tan\varphi_t. \quad (7)$$

## 2 计算模型

深埋洞室在开挖过程中, 由于支护不及时或强度不够都极易发生垮塌破坏. 基于已有研究成果<sup>[7,9,22]</sup>,

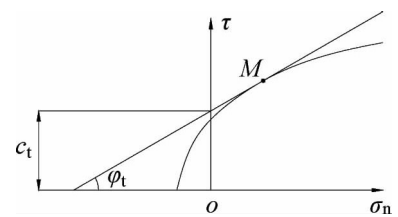


图1 非线性破坏准则与切线法

根据极限分析上限定理的要求,构建了深埋洞室的破坏模式以及对应的速度场,如图2所示.即洞室宽为  $l$ ,高为  $h$ ,洞室正上方楔形体  $ABGOC_1B_1$  以  $v_0$  的速度竖直向下垮塌,扇形体  $BGC_1$  与  $B_1G_1D_1$  分别绕  $G$  和  $G_1$  点发生转动破坏,三角形  $C_1GC_2, C_2GC_3, \dots, C_nGF$  以及  $D_1G_1D_2, D_2G_1D_3, \dots, D_nG_1F_1$  发生平动破坏.拱顶以及边墙产生的围岩压力分别为  $q, e$ ,且  $e = Kq$ .

根据图2b得,深埋洞室计算模型中各块体的速度分别为

$$v_1 = v_0; \quad (8)$$

$$v_2 = v_1 e^{\alpha_1 \tan 2\varphi_t} = v_0 e^{\alpha_1 \tan 2\varphi_t}; \quad (9)$$

$$v_{k,k+1} = \frac{\sin \alpha_k}{\cos 2\varphi_t} v_k \quad (k \in Z \quad 2 \leq k \leq n); \quad (10)$$

$$v_{k+1} = \frac{\cos(2\varphi_t - \alpha_k)}{\cos 2\varphi_t} v_k \quad (k \in Z \quad 2 \leq k \leq n). \quad (11)$$

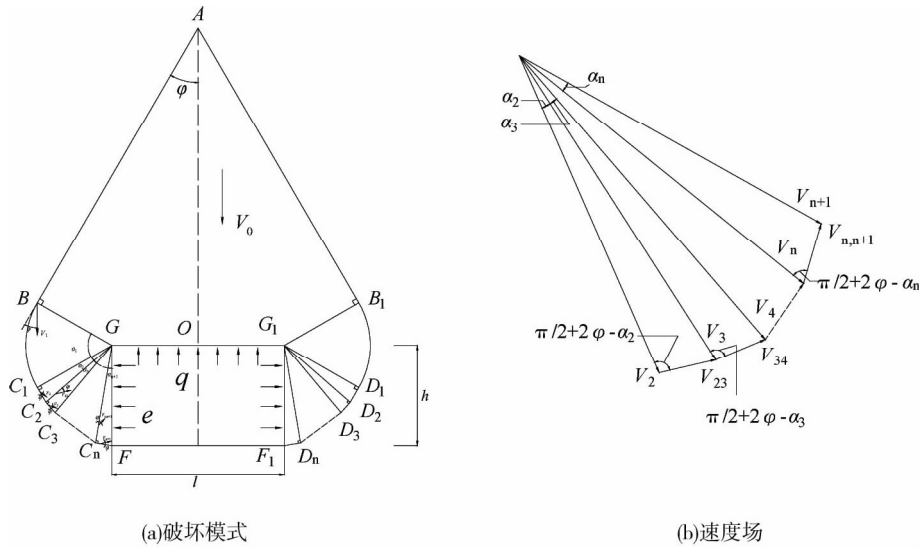


图2 深埋洞室的计算模型

### 3 推导过程

根据极限分析上限定理,可得到深埋洞室发生垮塌破坏时的外功率与内能耗散功率.

#### 3.1 外功率

##### 3.1.1 重力功率

$$W_{\text{soil}} = \gamma S_{ABCO} v_0 + \gamma \int_0^{\alpha_1} \frac{1}{2} GB^2 v_0 e^{\theta' \tan 2\varphi_t} \sin\left(\frac{\pi}{2} - \theta'\right) d\theta' + \sum_1^{n-1} \gamma S_{GC_i C_{i+1}} v_{i+1} \cos\left(\frac{\pi}{2} + \varphi_t - \sum_{i+1}^{n+1} \alpha_j\right) + \gamma S_{GC_n F} v_{n+1} \cos\left(\frac{\pi}{2} + \varphi_t - \alpha_{n+1}\right); \quad (12)$$

$$W' = \gamma S_{ABCO} v_0 = \gamma \left( \frac{1}{2} BGAB + \frac{1}{2} AOGO \right) v_0 = \gamma v_0 h^2 f'; \quad (13)$$

$$f' = \frac{1}{2} \prod_2^{n+1} \cos \alpha_i \left( \frac{\prod_2^{n+1} \cos \alpha_i}{\tan \varphi_t} + \frac{l}{2h \sin \varphi_t} \right) + \left[ \left( \frac{\prod_2^{n+1} \cos \alpha_i}{\tan \varphi_t} + \frac{l}{2h \sin \varphi_t} \right) \cos \varphi_t + \prod_2^{n+1} \cos \alpha_i \sin \varphi_t \right] \frac{l}{4h}; \quad (14)$$

$$W'' = \gamma \int_0^{\alpha_1} \frac{1}{2} GB^2 v_0 e^{\theta' \tan 2\varphi_t} \sin\left(\frac{\pi}{2} - \theta'\right) d\theta' = \gamma v_0 h^2 f''; \quad (15)$$

$$f'' = \frac{1}{2} \frac{\prod_2^{n+1} \cos^2 \alpha_i}{1 + \tan^2 2\varphi_t} \left[ (\sin \alpha_1 + \tan 2\varphi_t \cos \alpha_1) e^{\alpha_1 \tan 2\varphi_t} - \tan 2\varphi_t \right]; \quad (16)$$

$$W_k = \gamma S_{GC_k C_{k+1}} v_{k+1} \cos\left(\frac{\pi}{2} + \varphi_t - \sum_{k+1}^{n+1} \alpha_i\right) = \gamma v_0 h^2 f_k; \quad (17)$$

$$f_1 = \frac{1}{2} \sin \alpha_2 \cos \alpha_2 \left( \prod_3^{n+1} \cos^2 \alpha_i \right) \cos \left( \frac{\pi}{2} + \varphi_t - \sum_2^{n+1} \alpha_j \right) e^{\alpha_1 \tan 2\varphi_t}; \quad (18)$$

$$f_k = \frac{\prod_2^k \cos(2\varphi_t - \alpha_i)}{2 \cos^{k-1} 2\varphi_t} \sin \alpha_{k+1} \cos \alpha_{k+1} \left( \prod_{k+2}^{n+1} \cos^2 \alpha_i \right) \cos \left( \frac{\pi}{2} + \varphi_t - \sum_{k+1}^{n+1} \alpha_j \right) e^{\alpha_1 \tan 2\varphi_t};$$

$$(2 \leq k \leq n-1) \quad (19)$$

$$W_n = \gamma S_{GC_n F} v_{n+1} \cos \left( \frac{\pi}{2} + \varphi_t - \alpha_{n+1} \right) = \gamma v_0 h^2 f_n; \quad (20)$$

$$f_n = \frac{1}{2} \frac{\prod_2^n \cos(2\varphi_t - \alpha_i)}{\cos^{n-1} 2\varphi_t} \sin \alpha_{n+1} \cos \alpha_{n+1} \cos \left( \frac{\pi}{2} + \varphi_t - \alpha_{n+1} \right) e^{\alpha_1 \tan 2\varphi_t}. \quad (21)$$

### 3.1.2 支护反力功率

$$W_T = -q \frac{l}{2} v_0 - e h v_{n+1} \cos(\alpha_{n+1} - \varphi_t) = -q h v_0 f_0; \quad (22)$$

$$e = Kq; \quad (23)$$

$$f_0 = \frac{l}{2h} + K \frac{\prod_2^n \cos(2\varphi_t - \alpha_i)}{\cos^{n-2} 2\varphi_t} \cos(\alpha_{n+1} - \varphi_t) e^{\alpha_1 \tan 2\varphi_t}. \quad (24)$$

### 3.1.3 外力总功率

外力总功率等于围岩自重功率与支护反力功率之和,即:

$$W_{\text{ext}} = W_{\text{soil}} + W_T = \gamma v_0 h^2 (f' + f'' + \sum_1^n f_k) - q h v_0 f_0. \quad (25)$$

## 3.2 内能耗散功率

沿 AB 间断线的能量耗散:

$$D_{AB} = c_t A B v_0 \cos \varphi_t = c_t h v_0 g'; \quad (26)$$

$$g' = \left( \frac{\prod_2^{n+1} \cos \alpha_i}{\tan \varphi_t} + \frac{l}{2h \sin \varphi_t} \right) \cos \varphi_t. \quad (27)$$

圆弧受剪面  $BC_1$  和受剪区  $GBC_1$  的能量耗散:

$$D_{BC_1} = \frac{c_t B G v_1 \cos \varphi_t}{\tan 2\varphi_t} (e^{\alpha_1 \tan 2\varphi_t} - 1) = c_t h v_0 g''; \quad (28)$$

$$g'' = \frac{\prod_2^{n+1} \cos \alpha_i \cos \varphi_t}{\tan 2\varphi_t} (e^{\alpha_1 \tan 2\varphi_t} - 1); \quad (29)$$

$$D_{GBC} = \frac{c_t B G v_1 \cos \varphi_t}{\sin 2\varphi_t} (e^{\alpha_1 \tan 2\varphi_t} - 1) = c_t h v_0 g'''; \quad (30)$$

$$g''' = \frac{\prod_2^{n+1} \cos \alpha_i \cos \varphi_t}{\sin 2\varphi_t} (e^{\alpha_1 \tan 2\varphi_t} - 1). \quad (31)$$

沿间断线  $C_k C_{k+1}$  ( $k \in Z, 1 \leq k \leq n-1$ ) 的能量耗散:

$$D_{C_k C_{k+1}} = c_t C_k C_{k+1} v_{k+1} \cos \varphi_t = c_t h v_0 g_k; \quad (31)$$

$$g_1 = \prod_2^{n+1} \sin \alpha_i \cos \varphi_t e^{\alpha_1 \tan 2\varphi_t}; \quad (32)$$

$$g_k = \frac{\prod_2^k \cos(2\varphi_t - \alpha_i) \prod_{k+1}^{n+1} \sin \alpha_i}{\cos^{k-1}(2\varphi_t)} \cos \varphi_t e^{\alpha_1 \tan 2\varphi_t} (2 \leq k \leq n-1). \quad (33)$$

沿间断线  $C_n F$  的能量耗散:

$$D_{C_n F} = c_1 C_n F v_{n+1} \cos \varphi_t = c_1 h v_0 g_n; \quad (34)$$

$$g_n = \frac{\prod_2^n \cos(2\varphi_t - \alpha_i) \sin \alpha_{n+1}}{\cos^{n-1}(2\varphi_t)} \cos \varphi_t e^{\alpha_1 \tan 2\varphi_t}. \quad (35)$$

沿  $GC_k (k \in Z, 2 \leq k \leq n)$  间断线的能量耗散:

$$D_{GC_k} = c_1 GC_k v_{k,k+1} \cos \varphi_t = c_1 v_0 h l_k; \quad (36)$$

$$l_2 = \frac{\sin \alpha_2 \prod_3^{n+1} \cos \alpha_j}{\cos(2\varphi_t)} \cos \varphi_t e^{\alpha_1 \tan 2\varphi_t}; \quad (37)$$

$$l_k = \frac{\prod_2^{k-1} \cos(2\varphi_t - \alpha_i) \sin \alpha_k \prod_{k+1}^{n+1} \cos \alpha_j}{\cos^{k-1}(2\varphi_t)} \cos \varphi_t e^{\alpha_1 \tan 2\varphi_t} (3 \leq k \leq n). \quad (38)$$

总内能耗散功率:

$$D_{\text{int}} = D_{AB} + D_{BC_1} + D_{BC_2} + \sum_1^{n-1} D_{C_k C_{k+1}} + \sum_2^n D_{GC_k} + D_{C_n F} = c_1 h v_0 (g' + g'' + g''' + \sum_1^n g_k + \sum_2^n l_k). \quad (39)$$

### 3.3 围岩压力

根据虚功率原理,外功率与内能耗散功率相等,即:

$$W_{\text{ext}} = D_{\text{int}}. \quad (40)$$

由此可得,深埋洞室拱顶围岩压力的表达式为

$$q = \frac{\gamma h (f' + f'' + \sum_1^n f_k) - c_1 (g' + g'' + g''' + \sum_1^n g_k + \sum_2^n l_k)}{f_0}. \quad (41)$$

## 4 结果分析

根据文献[22],在三角形块体  $n=3$  时可以满足工程实践的精度要求.

当  $n=3$  时,非线性 Mohr-Coulomb 破坏准则下深埋洞室围岩压力的表达式为

$$q = \frac{\gamma h (f_1 + f_2 + f_3 + f_4 + f_5) - c_1 (f_7 + f_8 + f_9 + f_{10} + f_{11} + f_{12} + f_{13} + f_{14})}{f_6}. \quad (42)$$

式中:

$$c_1 = \frac{m-1}{m} c_0 \left( \frac{m \sigma_t \tan \varphi_t}{c_0} \right)^{\frac{1}{1-m}} + \sigma_t \tan \varphi_t; \quad (43)$$

$$f_1 = \frac{1}{2} \cos \alpha_2 \cos \alpha_3 \cos \alpha_4 \left( \frac{\cos \alpha_2 \cos \alpha_3 \cos \alpha_4}{\tan \varphi_t} + \frac{l}{2h \sin \varphi_t} \right) + \left[ \left( \frac{\cos \alpha_2 \cos \alpha_3 \cos \alpha_4}{\tan \varphi_t} + \frac{l}{2h \sin \varphi_t} \right) \cos \varphi_t + \cos \alpha_2 \cos \alpha_3 \cos \alpha_4 \sin \varphi_t \right] \frac{l}{4h}; \quad (44)$$

$$f_2 = \frac{1}{2} \frac{\cos^2 \alpha_2 \cos^2 \alpha_3 \cos^2 \alpha_4}{1 + \tan^2 2\varphi_t} [(\sin \alpha_1 + \tan 2\varphi_t \cos \alpha_1) e^{\alpha_1 \tan 2\varphi_t} - \tan 2\varphi_t]; \quad (45)$$

$$f_3 = \frac{1}{2} \sin \alpha_2 \cos \alpha_2 \cos^2 \alpha_3 \cos^2 \alpha_4 \sin(\alpha_2 + \alpha_3 + \alpha_4 - \varphi_t) e^{\alpha_1 \tan 2\varphi_t}; \quad (46)$$

$$f_4 = \frac{1}{2} \sin \alpha_3 \cos \alpha_3 \cos^2 \alpha_4 \frac{\cos(2\varphi_t - \alpha_2)}{\cos 2\varphi_t} e^{\alpha_1 \tan 2\varphi_t} \sin(\alpha_3 + \alpha_4 - \varphi_t); \quad (47)$$

$$f_5 = \frac{1}{2} \sin \alpha_4 \cos \alpha_4 \frac{\cos(2\varphi_t - \alpha_3) \cos(2\varphi_t - \alpha_2)}{\cos^2 2\varphi_t} e^{\alpha_1 \tan 2\varphi_t} \sin(\alpha_4 - \varphi_t); \quad (48)$$

$$f_6 = \frac{l}{2h} + K \frac{\cos(2\varphi_t - \alpha_3) \cos(2\varphi_t - \alpha_2)}{\cos^2 2\varphi_t} e^{\alpha_1 \tan 2\varphi_t} \cos(\alpha_4 - \varphi_t); \quad (49)$$

$$f_7 = \left( \frac{\cos\alpha_2 \cos\alpha_3 \cos\alpha_4}{\tan\varphi_t} + \frac{l}{2h\sin\varphi_t} \right) \cos\varphi_t; \quad (50)$$

$$f_8 = \frac{\cos\alpha_4 \cos\alpha_3 \cos\alpha_2 \cos\varphi_t}{\tan 2\varphi_t} (e^{\alpha_1 \tan 2\varphi_t} - 1); \quad (51)$$

$$f_9 = \frac{\cos\alpha_4 \cos\alpha_3 \cos\alpha_2 \cos\varphi_t}{\sin 2\varphi_t} (e^{\alpha_1 \tan 2\varphi_t} - 1); \quad (52)$$

$$f_{10} = \sin\alpha_2 \cos\alpha_3 \cos\alpha_4 e^{\alpha_1 \tan 2\varphi_t} \cos\varphi_t; \quad (53)$$

$$f_{11} = \frac{\sin\alpha_2 \cos\alpha_3 \cos\alpha_4 \cos\varphi_t}{\cos 2\varphi_t} e^{\alpha_1 \tan 2\varphi_t}; \quad (54)$$

$$f_{12} = \frac{\cos(2\varphi_t - \alpha_2) \sin\alpha_3 \cos\alpha_4 \cos\varphi_t}{\cos 2\varphi_t} e^{\alpha_1 \tan 2\varphi_t}; \quad (55)$$

$$f_{13} = \frac{\cos\varphi_t \cos\alpha_4 \sin\alpha_3 \cos(2\varphi_t - \alpha_2)}{\cos^2 2\varphi_t} e^{\alpha_1 \tan 2\varphi_t}; \quad (56)$$

$$f_{14} = \frac{\sin\alpha_4 \cos\varphi_t \cos(2\varphi_t - \alpha_3) \cos(2\varphi_t - \alpha_2)}{\cos^2 2\varphi_t} e^{\alpha_1 \tan 2\varphi_t}. \quad (57)$$

角度参数的约束条件为

$$\alpha_1 + \alpha_2 + \alpha_3 + \alpha_4 = \pi/2 + \varphi_t. \quad (58)$$

根据文献[22]的计算过程,得到无数个 $q$ 值,其中最大的 $q$ 值即为围岩压力的上限解.

#### 4.1 非线性系数与侧压力比例系数对围岩压力的影响

取各参数:土体重度 $\gamma = 20 \text{ kN/m}^3$ , 硐室宽度 $l = 10 \text{ m}$ , 高度 $h = 10 \text{ m}$ , 初始粘聚力 $c_0 = 10 \text{ kPa}$ , 单轴抗拉强度 $\sigma_t = 50 \text{ kPa}$ ; 非线性系数 $m$ 分别取1.1, 1.2, 1.3, 1.4, 1.5, 1.6, 侧压力比例系数 $K$ 分别取0.4, 0.6, 0.8, 1.0, 1.2, 1.4. 根据所构建的破坏机制,采用极限分析上限法所得到的围岩压力的上限解如图3所示.由图3可得,随着非线性系数 $m$ 的增大,拱顶以及边墙的围岩压力 $q, e$ 均呈非线性增大,当非线性系数较小时,增加的幅度较小;而当非线性系数较大时,增大的效果非常明显,由此可见非线性系数是影响深埋硐室拱顶以及边墙围岩压力的主要因素之一.随着侧压力比例系数 $K$ 的增大,拱顶围岩压力 $q$ 减小,边墙围岩压力 $e$ 增大,这说明了侧压力比例系数的变化显著影响着拱顶、边墙围岩压力的大小,此结论与杨小礼等<sup>[23]</sup>、杨峰等<sup>[24]</sup>研究成果一致.拱顶的支护设计固然重要,但是当侧压力比例系数变大时,边墙承受的围岩压力增大,此时可理解为边墙分担了拱顶一部分的围岩压力,故拱顶承受的围岩压力随之减小,因此,边墙的支护设计对隧道整体的稳定性具有重要的意义,不容忽视.

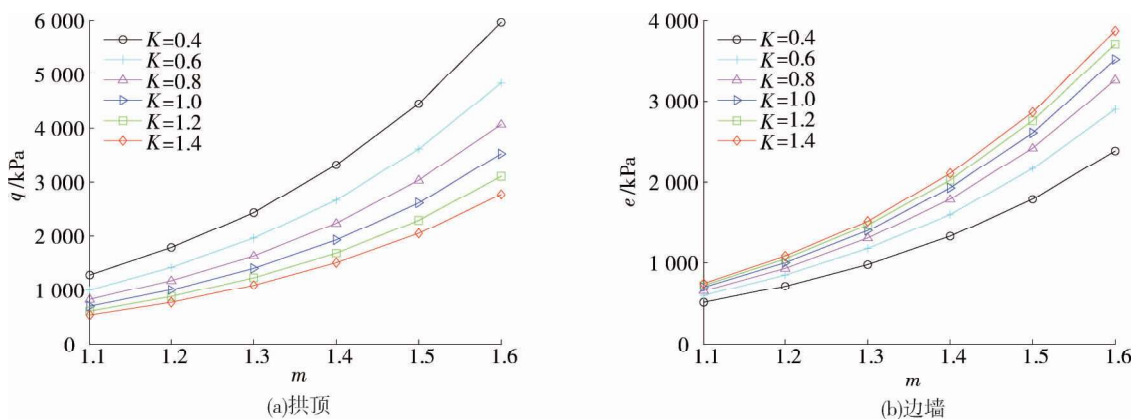


图3 非线性系数与侧压力比例系数对围岩压力的影响

#### 4.2 硐室断面尺寸对围岩压力的影响

研究硐室断面尺寸对围岩压力的影响时,取各参数为:土体重度 $\gamma = 20 \text{ kN/m}^3$ , 非线性系数 $m = 1.1$ , 初始粘聚力 $c_0 = 10 \text{ kPa}$ , 单轴抗拉强度 $\sigma_t = 50 \text{ kPa}$ , 侧压力比例系数 $K = 1.0$ ; 硐室宽度 $l$ 和高度 $h$ 取5, 6, 7, 8, 9, 10 m. 根据极限分析上限法所求得的围岩压力上限解如图4所示,当侧压力比例系数 $K = 1.0$ 时,拱

顶以及边墙的围岩压力是相等的,即  $q = e$ 。随着洞室断面尺寸的增大,拱顶以及边墙的围岩压力均呈增大的趋势.这说明了洞室断面尺寸影响着拱顶、边墙围岩压力的大小,当洞室断面尺寸较小时,拱顶和边墙的围岩压力较小,而当洞室断面尺寸较大时,拱顶和边墙的围岩压力较大,此时应根据拱顶、边墙围岩压力的大小进行支护设计,选择合理经济的支护方案。

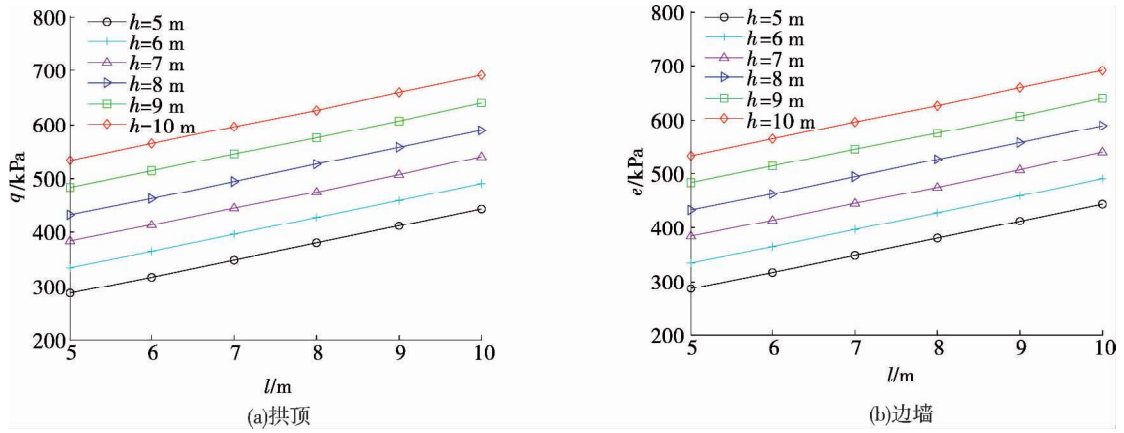


图 4 洞室断面尺寸对围岩压力的影响

### 4.3 初始粘聚力与轴向拉应力对围岩压力的影响

取各参数为:土体重度  $\gamma = 20 \text{ kN/m}^3$ , 非线性系数  $m = 1.1$ , 洞室宽度  $l = 10 \text{ m}$ , 高度  $h = 10 \text{ m}$ , 侧压力比例系数  $K = 1.0$ ; 初始粘聚力  $c_0$  分别取 5, 10, 15, 20, 25, 30 kPa, 轴向拉应力  $\sigma_i$  分别取 50, 60, 70, 80, 90, 100 kPa. 根据极限分析上限法所求得的围岩压力上限解如图 5 所示. 当侧压力比例系数  $K = 1.0$  时, 拱顶以及边墙的围岩压力是相等的, 即  $q = e$ 。随着初始粘聚力  $c_0$  增大, 拱顶以及边墙的围岩压力  $q, e$  非线性减小, 当初始粘聚力较小时, 减小的趋势较陡, 而当初始粘聚力较大时, 减小的趋势较为平缓. 可见初始粘聚力的持续减小会引起拱顶以及边墙围岩压力的陡然增大, 尤其针对软弱围岩, 更应该引起重视, 需加强拱顶以及边墙的支护措施. 随着轴向拉应力  $\sigma_i$  的增大, 拱顶以及边墙的围岩压力  $q, e$  增大, 但效果不如初始粘聚力明显。

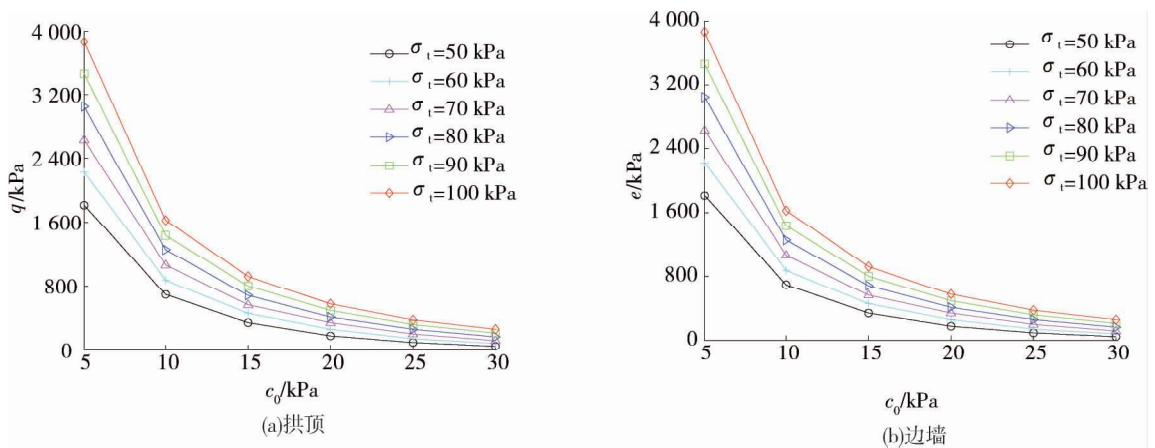


图 5 初始粘聚力与轴向拉应力对围岩压力的影响

## 5 结论

1) 在非线性 Mohr-Coulomb 破坏准则下, 随着非线性系数的增大, 深埋洞室拱顶以及边墙的围岩压力非线性增大, 且效果非常显著. 因此非线性系数是影响深埋洞室围岩压力的主要因素, 同时也说明了采用非线性 Mohr-Coulomb 破坏准则计算深埋洞室的围岩压力是非常有必要的。

2) 洞室断面尺寸、初始粘聚力以及轴向拉应力在不同程度上影响着深埋洞室的围岩压力。

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